

2000-10-15

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Itay Furman, Ofer Biham, Jiang-Kai Zuo, Anna K Swan, John F Wendelken. 2000. "Epitaxial growth of Cu on Cu(001): Experiments and simulations." *Physical Review B*, Volume 62, Issue 16, R10652(R). <https://doi.org/10.1103/physrevb.62.r10649>

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Epitaxial growth of Cu on Cu(001): Experiments and simulations

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(Received 31 July 2000)

A quantitative comparison between experimental and Monte Carlo simulation results for the epitaxial growth of Cu/Cu(001) in the submonolayer regime is presented. The simulations take into account a complete set of hopping processes whose activation energies are derived from semiempirical calculations using the embedded-atom method. The island separation is measured as a function of the incoming flux and the temperature. A good *quantitative* agreement between the experiment and simulation is found for the island separation, the activation energies for the dominant processes, and the exponents that characterize the growth. The simulation results are then analyzed at lower coverages, which are not accessible experimentally, providing good agreement with theoretical predictions as well.

The growth of a deposited layer (adlayer) of a thin metal film in molecular-beam epitaxy (MBE) through island nucleation involves three kinetic regimes with respect to the coverage θ : (a) *nucleation regime*, dominated by island nucleation; (b) *aggregation regime*, where deposited atoms (adatoms) are mostly captured by existing islands. In this regime the system is in a quasi steady-state, a property that is employed in the scaling theory; (c) *coalescence regime*, where islands are large enough to interact and merge.

Experiments using scanning tunneling microscopy (STM),¹⁻³ low-energy electron diffraction⁴⁻⁷ and other methods,⁸ made in attempt to study the aggregation regime, reveal an exponential dependence of the island density N , on the inverse temperature $1/T$,^{1,2,5,7,8} as well as a power law dependence on the flux F .^{2-4,6,8} This behavior can be described by the following formula:⁹

$$\ell^{-2} \sim \left(\frac{F}{\nu}\right)^{\gamma} \exp(E_{\text{eff}}/k_B T), \quad (1)$$

where ℓ is the mean island separation related to N through $N \sim \ell^{-2}$. In the experimental-data analysis, γ , ν , and E_{eff} are fitting parameters, where E_{eff} is a measure of the activation energy for the rate-limiting move, ν is the attempt rate for that move and γ is a scaling exponent. The theoretical challenge is, therefore, to relate γ , ν , and E_{eff} to the fundamental processes at the atomic level.

In this paper we show that kinetic Monte Carlo (MC) simulations based on energy barriers obtained via the embedded-atom method (EAM), provide a good quantitative description of diffusion and growth on Cu(001). To this end we performed kinetic MC simulations of island growth under conditions identical to those employed experimentally in Refs. 4–6. In these experiments spot-profile analysis of low-energy electron diffraction (SPA-LEED) was used to measure the mean separation ℓ between Cu islands on Cu(001). The measurements were taken as a function of the flux F for three temperatures, $T=213$, 223, and 263 K (Fig. 1), and as

a function of the temperature T at a constant flux $F=3.21 \times 10^{-4} \text{ ML s}^{-1}$ (Fig. 2).

In our kinetic MC simulations atoms are randomly deposited on a square-lattice substrate of 250×250 sites that corresponds to a terrace width of $\sim 64 \text{ nm}$.¹⁰ These atoms attach irreversibly to the surface and hop as random walkers to unoccupied nearest-neighbor (NN) sites. Each hop involves an activation energy-barrier, E_n , that depends on the configuration of occupied and unoccupied adjacent sites in the 3×3 square around the hopping atom.¹¹⁻¹³ When an atom is deposited on top of an island it is incorporated at a random position along the island perimeter. The nucleation of a second layer is thus suppressed, which is a good approximation for small island sizes at low coverages. The hopping-rate (in units of hops per second), h_n , for some configuration n is¹⁴

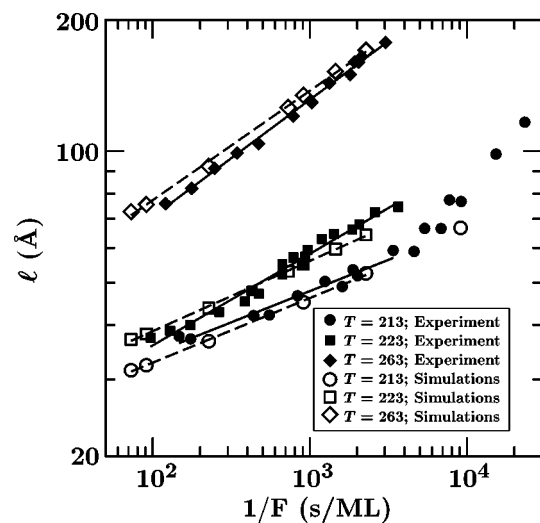


FIG. 1. Comparison of experimental results (full symbols) and simulation results (empty symbols) for the island separation ℓ vs the inverse flux $1/F$, at three temperatures $T=213$, 223, and 263 K. The coverage is $\theta=0.3 \text{ ML}$. The solid lines represent fits to Eq. (1). The resultant values of the exponent γ are presented in Table I.

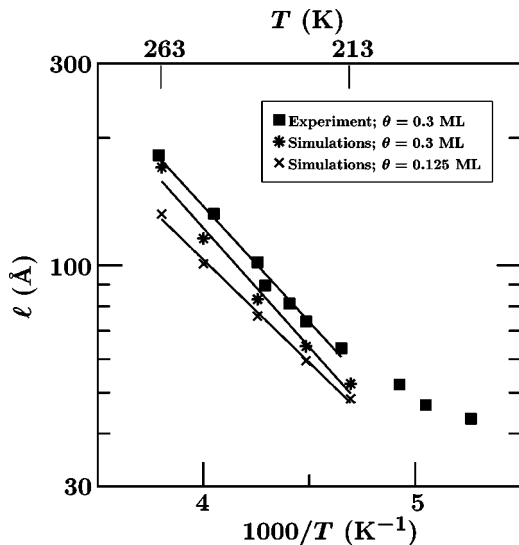


FIG. 2. The island separation l vs $1/T$, obtained from the experiment for $F=3.21 \times 10^{-4} \text{ ML s}^{-1}$ and simulation for $F=4.8 \times 10^{-4} \text{ ML s}^{-1}$ at coverage $\theta=0.3 \text{ ML}$. We also present the simulation results of the same runs but at an earlier stage when $\theta=0.125 \text{ ML}$. The solid lines represent fits to $l \sim \exp(-E_{\text{eff}}/2K_B T)$, with $E_{\text{eff}}=0.108 \pm 0.005$, 0.112 ± 0.008 , and $0.096 \pm 0.003 \text{ eV}$ for the experiment, and the simulations at $\theta=0.3$ and 0.125 ML , respectively.

$$h_n = \nu \exp(-E_n/K_B T), \quad (2)$$

where ν is the attempt rate common to all moves and E_n is calculated using the EAM.^{15,16,13}

We have reconstructed the experimental conditions in our kinetic MC simulations and obtained for each experimental curve in Figs. 1 and 2 a corresponding simulated curve. Each simulated data point is an average over 20 runs. In Fig. 1 we present the experimental results and the corresponding simulation results, for the island separation l vs $1/F$ at $T=213$, 223 , and 263 K . All data points were taken at $\theta=0.3 \text{ ML}$.¹⁷ Clearly, there is a good agreement between the experiment and simulation. The curves follow a power-law behavior according to Eq. (1) in a range of one and a half decades, from which the exponent γ is extracted.¹⁸ The results for the exponent γ are presented in Table I.¹⁹ It is found to increase in the range $\gamma=0.3-0.5$ as the temperature is raised, suggesting that the system undergoes a crossover in its kinetic behavior.²⁰ Based on scaling theory, one may conclude from the results $\gamma \approx 1/3$ that *dimers* and larger islands are stable and *immobile*, and adatoms are the *only* mobile entities on the surface.⁹ However, care should be exercised when draw-

TABLE I. The values for the exponent γ defined in Eq. (1), for $\theta=0.3 \text{ ML}$ (Fig. 1) and $\theta=0.125 \text{ ML}$.

$T \text{ (K)}$	γ		
	Experiment $\theta=0.3 \text{ (ML)}$	Simulations $\theta=0.3 \text{ (ML)}$	Simulations $\theta=0.125 \text{ (ML)}$
213	0.28 ± 0.02	0.30 ± 0.01	0.309 ± 0.006
223	0.42 ± 0.02	0.326 ± 0.006	0.319 ± 0.006
263	0.54 ± 0.01	0.51 ± 0.01	0.396 ± 0.006

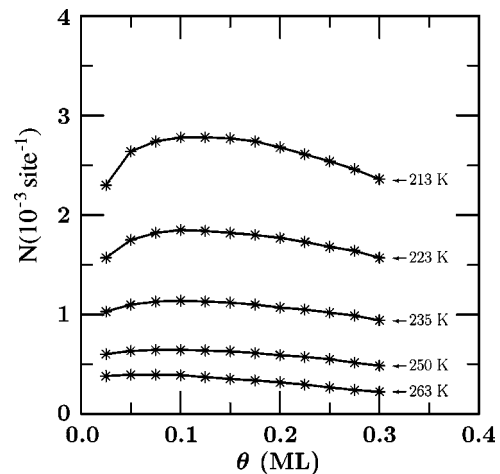


FIG. 3. The island density N , vs the coverage θ at different temperatures with a flux $F=4.8 \times 10^{-4} \text{ ML s}^{-1}$. The solid lines are a guide to the eye. Since all the runs begin with an empty surface, the point $(0, 0)$ is common to all the curves although it is not drawn in the figure.

ing such conclusions. Specifically, in our case the value $\gamma \approx 1/3$ is found *in spite of significant dimer mobility* as we show later.

Unlike the energy barriers, we do not have a value for the attempt rate ν , from atomic scale calculations. Instead, the *experimental* value of ν is obtained by fitting the simulation and experimental results. This is possible since in the simulation ν sets a fundamental clock rate, while the simulation results depend only on the ratio ν/F rather than on F and ν separately. Using this property we perform the simulations for a broad range of values of ν/F and plot the island separation l vs ν/F . The attempt rate ν , is then obtained as the value for which the three simulated curves in Fig. 1, overlap simultaneously the three experimental curves. It is found to be $\nu=1.2 \times 10^{13} \text{ s}^{-1}$.

In Fig. 2 we present the experimental results (■) for the island separation l vs $1/T$, at $F=3.21 \times 10^{-4} \text{ ML s}^{-1}$ and the corresponding simulation results (*),²¹ taken at $\theta=0.3 \text{ ML}$. The apparent good agreement between them is a further confirmation of the correctness of ν . Both experimental and simulation curves follow the Arrhenius behavior predicted by Eq. (1). In these plots the slope of the curve corresponds to the activation energy E_{eff} for the rate limiting step of the nucleation process. The results are $E_{\text{eff}}=0.108 \pm 0.005$ and $0.112 \pm 0.008 \text{ eV}$ for the experimental and simulation curves, respectively.

In Fig. 3 we present the simulation results for the island density N vs the coverage θ , for $F=4.8 \times 10^{-4} \text{ ML s}^{-1}$ and different temperatures within the experimental range $213-263 \text{ K}$. The data extracted from the simulations is marked by (*), while the connecting lines are only guides to the eye. The quick rise of N in the nucleation regime, is followed by a plateau in the aggregation regime, and then a slow decrease that is the mark of the coalescence regime. The coalescence at the lower temperatures $T=213$ and 223 K , is system-size independent and is due to island separation being comparable to island size. At the higher temperatures, especially $T=263 \text{ K}$, the coalescence is mainly due to island separation being comparable to system size, i.e., finite size effects. We believe that this effect reproduces to a certain extent the ef-

TABLE II. The adatom diffusion coefficient and the relative dimer and trimer diffusion coefficients for the three studied temperatures. Note that through the whole temperature range the relation between the diffusion coefficients almost does not change. The diffusion coefficients of islands of size four and higher are orders of magnitude lower.

T (K)	D_1 (sites s^{-1})	D_2/D_1	D_3/D_1
213	3.6 ± 0.1	1.5 ± 0.1	0.021 ± 0.002
223	11.7 ± 0.4	1.3 ± 0.1	0.023 ± 0.002
263	520.0 ± 15.0	1.27 ± 0.08	0.040 ± 0.003

fect of the steps limiting the width of the terraces in the experimental system.¹⁰ Our simulations indicate that at $\theta = 0.3$ ML, where the experiments were done, the system is already in the coalescence regime. Therefore, the comparison of both experiment and simulations with theory [Eq. (1)] is misleading.

The comparison between the experimental and simulation results shows that the set of energy barriers used here, provides a good quantitative description of the submonolayer growth in the low coverage regime. We now proceed to analyze the results in the context of available theoretical work. In Table II we present the calculated diffusion coefficients of adatoms as well as of small islands (dimers and trimers) at the experimentally relevant temperatures. The diffusion coefficient of a dimer D_2 , is found to be comparable to the adatom diffusion coefficient D_1 . The diffusion coefficients of trimers and larger islands (not shown) are almost two orders of magnitude smaller. Thus, only adatoms and dimers are mobile on the surface on a time scale relevant to the growth process. In this case the mean island separation in the aggregation regime is given by:^{12,13,22}

$$l \sim (D_1 D_2 / F^2)^{1/10}. \quad (3)$$

In Ref. 13 it was shown that the mobility of monomers and dimers is fully determined by only three energy barriers, the single adatom hopping E_0 , the dimer lateral-bond breaking E_2 and the re-establishing of a NN bond E_4 (the labeling follows Refs. 13 and 16):

$$E_0 = E \left(\ominus \right) \quad E_2 = E \left(\text{⊗} \right) \quad E_4 = E \left(\text{⊗} \right).$$

To determine also the dimer stability one needs to specify the dimer bond-breaking energy E_8 , and the next-nearest-neighbor bond-breaking energy E_1 :

$$E_1 = E \left(\text{⊙} \right) \quad E_8 = E \left(\text{⊗} \right).$$

Specifically, in the case of Cu, $E_0 = 0.485$, $E_1 = 0.563$ (In Ref. 16 the value of this barrier was mistyped.), $E_2 = 0.463$, $E_4 = 0.183$, and $E_8 = 0.811$ eV. Dimer breaking can take place either from the NN configuration with energy barrier E_8 , or from the next-nearest-neighbor configuration with effective energy barrier $E_2 - E_4 + E_1 = 0.843$ eV. The high values of these barriers guarantee the dimer stability. The adatom diffusion is a single step process, hence, $D_1 = \nu \exp(-E_0/K_B T)$ [Eq. (2)]. The dimer diffusion is a

double-step process involving a move with barrier E_2 , followed by a move with E_4 . The former move is the rate-limiting move, thus, up to a combinatorial factor of order 1 we approximate $D_2 = \nu \exp(-E_2/K_B T)$. Since $E_2 \approx E_0$ we expect D_2 and D_1 to be comparable, as is indeed the case (see Table II). Note that this energy structure is predicted to be common to most of the fcc(001) metal surfaces.¹⁶

Inserting D_1 and D_2 into Eq. (3) and comparing with Eq. (1) we get

$$E_{\text{eff}} = (E_0 + E_2)/10 = 0.0948 \text{ eV}, \quad (4)$$

that differs by more than 10% from the values obtained from the simulations ($E_{\text{eff}} = 0.112 \pm 0.008$ eV) and the experiment ($E_{\text{eff}} = 0.108 \pm 0.005$ eV) presented above in Fig. 2. We recall that the evaluation of E_{eff} was done at $\theta = 0.3$ ML where the system is already in the coalescence regime, and therefore deviations from Eqs. (1), (3), and (4) are expected. Therefore, to examine the scaling relations of Eqs. (3) and (4) we recalculate the island separation l vs $1/T$ in the same temperature range for a lower coverage of $\theta = 0.125$ ML where the validity of those equations is better satisfied. We find that $E_{\text{eff}} = 0.096 \pm 0.003$ eV is in good agreement with Eq. (4). These results are in excellent agreement with recent calculations by Boisvert and Lewis.²³

We will now revisit the calculation of γ . According to the scaling theory [Eq. (3)] $\gamma = 2/5$ for a system with mobile monomers and dimers.^{12,22} The experimental results for γ and the corresponding simulation results, obtained at $\theta = 0.3$ ML and presented in Table I are significantly different from that value. The discrepancy is removed when we recalculate γ for $T = 263$ K and a *much lower coverage*, $\theta = 0.125$ ML, where Eq. (3) applies. The results of this recalculation are presented in the rightmost column of Table I. Indeed we find $\gamma \approx 2/5$, in good agreement with the scaling prediction. For $T = 213$ and 223 K in the studied flux range the aggregation regime is wiped out, and the system is dominated by either nucleation or coalescence. Therefore, deviations from the scaling prediction are expected even for the lower coverage $\theta = 0.125$ ML, as observed in Table I.

The diffusion and coarsening of large Cu islands on Cu(001) has recently been studied experimentally.²⁴ The results were analyzed using MC simulation and a model that is qualitatively consistent with our model for the barriers relevant to large island diffusion.²⁵ However, the models differ significantly in the lateral-bond breaking barrier, E_2 , that is dominant in the monolayer growth through island nucleation. Actually, the crucial factor is the difference $E_2 - E_0 = 0.463 - 0.485 = -0.022$ eV in our model, and $0.52 - 0.399 = 0.121$ eV in the model of Heinonen *et al.*²⁵ It follows that while $D_2/D_1 = \exp[-(E_2 - E_0)/K_B T]$ is between 3.3 and 2.6 in the present model for temperatures between 213 and 263 K, it is between 0.001 and 0.005 in the model of Heinonen *et al.*²⁵ Thus, according to their model dimers are *static* entities on the surface giving rise to different scaling properties. On the other hand, the barriers of Shi *et al.*,²⁶ $E_0 = 0.503$ and $E_2 = 0.494$ eV, are in good agreement with ours.

In summary, we have presented a comparison between kinetic MC simulations and experiments of Cu/Cu(001) growth in the submonolayer regime. The available experi-

mental data is for $\theta=0.3$ ML and $T=213, 223,$ and 263 K. At this coverage a good agreement between simulation and experimental results is found. However, these results deviate from mean-field predictions, a fact that we attribute to the system being away from the aggregation regime at $\theta=0.3$ ML. Indeed, repeating the analysis of the simulation results for $\theta=0.125$ ML we find $\gamma\sim 0.4$ and $E_{\text{eff}}=0.096\approx(E_0+E_2)/10$ eV consistent with mean-field predictions when only adatoms and dimers are mobile on the surface. We conclude that for Cu/Cu(001), in the studied temperature range, only adatoms and dimers are mobile. Furthermore,

mean-field predictions apply at $\theta\sim 0.1$ ML but not at much higher coverage. Finally, the comparison between simulation and experiment, augmented with a scaling procedure, enables us to depict the attempt rate $\nu=1.2\times 10^{13}\text{ s}^{-1}$ of the experimental system.

We thank Hanoch Mehl for helpful discussions. Experimental portions of this work were supported by Oak Ridge National Laboratory, which is managed by Lockheed Martin Energy Research Corporation for the U.S. Department of Energy under Contract No. DE-AC05-96OR22464.

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¹⁷This value is above the minimum measurable coverage of $\theta=0.2$ ML. For the data obtained at higher coverage $\theta>0.6$ ML (Ref. 4) the suppression of second-layer nucleation in our simulations may not be appropriate. Also, scaling theories do not apply in this regime. Therefore, we concentrated only on the $\theta=0.3$ data.

¹⁸Note that the experimental curve at $T=223$ K has a slight S-shaped modulation (see also Fig. 3 in Ref. 4). This gives rise to the deviation in the exponent γ compared to the simulation results. However, the actual differences in the data points between the experiment and simulation are rather small.

¹⁹In the present analysis we have used the broadest possible range common to all data sets, that provides simultaneous high-quality fits for all of them. In Refs. 4 and 6, however, a narrower flux range was employed, giving rise to different values.

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